

## RESEARCH ARTICLE

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## Key Points:

- Formation of local minimums in phase space density profiles of ultrarelativistic electrons coincides with EMIC wave occurrence
- New local minimums are accompanied by narrowing of normalized pitch angle distributions
- The local minimums are indicative of EMIC wave-induced loss in the heart of the outer belt

## Supporting Information:

- Supporting Information S1

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## Signatures of Ultrarelativistic Electron Loss in the Heart of the Outer Radiation Belt Measured by Van Allen Probes

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**Abstract** Up until recently, signatures of the ultrarelativistic electron loss driven by electromagnetic ion cyclotron (EMIC) waves in the Earth's outer radiation belt have been limited to direct or indirect measurements of electron precipitation or the narrowing of normalized pitch angle distributions in the heart of the belt. In this study, we demonstrate additional observational evidence of ultrarelativistic electron loss that can be driven by resonant interaction with EMIC waves. We analyzed the profiles derived from Van Allen Probe particle data as a function of time and three adiabatic invariants between 9 October and 29 November 2012. New local minimums in the profiles are accompanied by the narrowing of normalized pitch angle distributions and ground-based detection of EMIC waves. Such a correlation may be indicative of ultrarelativistic electron precipitation into the Earth's atmosphere caused by resonance with EMIC waves.

### 1. Introduction

Wave-particle interactions causing loss and acceleration of electrons in the Earth's radiation belts have been extensively studied since the beginning of the space era (Imhof et al., 1977; Millan & Thorne, 2007; Shprits, Elkington, et al., 2008; Shprits, Subbotin, et al., 2008; Thorne, 2010; Thorne & Kennel, 1971; Xiao et al., 2009, 2010). Recently, particular attention has been paid to the dynamics of very energetic ultrarelativistic electrons (energy above  $\sim 1-2$  MeV) (e.g., Baker, Kanekal, Hoxie, Henderson, et al., 2013; Shprits et al., 2016; 2013; Xiao et al., 2015). However, major mechanisms controlling this population are still under debate. The purpose of this study is to provide additional observational evidence of ultrarelativistic electron loss due to local resonant interaction with electromagnetic ion cyclotron (EMIC) waves.

The launch of the Van Allen Probes, formerly known as Radiation Belt Storm Probes (RBSP), on 30 August 2012 has led to significant progress in understanding the dynamics of ultrarelativistic electron population. The Relativistic Electron-Proton Telescope (REPT) (Baker, Kanekal, Hoxie, Batiste, et al., 2013), which is a part of the Energetic Particle, Composition, and Thermal Plasma (ECT) Suite (Spence et al., 2013) on board the satellites, allows measuring the electron distribution of energies from  $\sim 2$  to  $\sim 10$  MeV and above and covers both inner and outer belt regions from  $1.2 R_E$  to  $5.8 R_E$ . A few days after the launch, the Van Allen Probes detected the unusual three-zone radiation belt structure (Baker, Kanekal, Hoxie, Henderson, et al., 2013). The third belt (narrow "storage ring") (Baker, Kanekal, Hoxie, Henderson, et al., 2013) persisting for almost a month was formed by ultrarelativistic electrons at  $L^* = \sim 3.25$  after the flux dropout on approximately 3 September 2012. Shprits et al. (2013) analyzed this unique event, modeling the electron dynamics with the Versatile Electron Radiation Belt (VERB) code (Shprits et al. 2009; Subbotin & Shprits, 2009). They showed that EMIC wave-induced electron scattering is responsible for the formation of the third belt, and the expanded plasmasphere placed the belt in a different plasma environment, where electromagnetic waves weakly affect ultrarelativistic electrons, and the third belt remained for tens of days. Long-term modeling results obtained by Drozdov et al. (2015) using the VERB code showed that observed 3.6 MeV electron fluxes are significantly overestimated if EMIC waves are not included in the simulations. Further comparison of the simulation results with the observations of the 17 January 2013 storm (Shprits et al., 2016) confirmed that EMIC waves cause fast local loss of the ultrarelativistic electrons in the outer belt.

Previous theoretical studies suggested that EMIC waves can effectively scatter sub-MeV and MeV electrons (Horne & Thorne, 1998; Summers & Thorne, 2003; Thorne & Kennel, 1971; Ukhorskiy et al., 2010). Supporting the theoretical findings, a number of observational case studies used VLF transmitter and receiver systems, riometers, balloons, and polar satellites to measure ultrarelativistic electron precipitation associated with simultaneous EMIC wave detection by ground-based or satellite magnetometers (Blum et al., 2015; Clilverd et al., 2007, 2015; Miyoshi et al., 2008; Rodger et al., 2008, 2015). Using a recently developed algorithm for determination of precipitation events of sub-MeV and MeV electrons from Polar Orbiting Environmental Satellites (POES) and Meteorological Operational (METOP-2) satellite (Carson et al., 2013), Hendry et al. (2016) showed that 60% to 90% of precipitation events coincide with the waves detected on the ground. Yet the observational studies of electron precipitation into the atmosphere suggested that sub-MeV and MeV electrons can be scattered by EMIC waves, such studies consider only the electron population inside the loss cone, leaving aside the effects of the waves on the pitch angle distribution of particles in the heart of the belt.

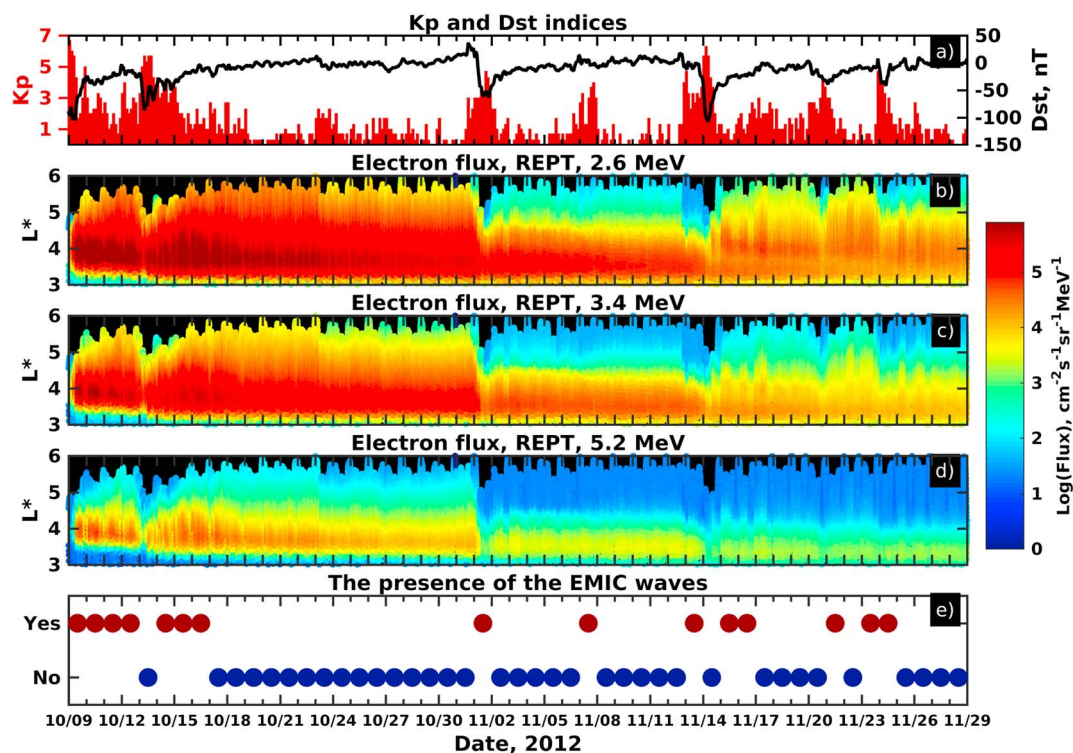
Usanova et al. (2014) investigated the correlation between ground-based EMIC wave observations from the Canadian Array for Real-time Investigations of Magnetic Activity (CARISMA) (Mann et al., 2008) and the variability of the ultrarelativistic outer belt measured in situ by the Van Allen Probes. They found the correlation between the observations of EMIC waves and narrowing of the normalized pitch angle distributions, which is a clear tell-tale signature of EMIC wave-induced precipitation.

Another mechanism of electron loss affecting all populations of radiation belt particles is the loss into the interplanetary medium driven by outward radial diffusion (Ohtani et al., 2009; Shprits et al., 2006; Turner et al., 2012). This mechanism is most effective during compression of the magnetosphere when particles previously trapped in the Earth's magnetic field at high radial distances find themselves drifting along the open trajectories that cross the magnetopause. These particles become lost into the interplanetary medium on time scales of the drift period, creating sharp negative gradients near the outer boundary of the outer belt. The sharp gradients, in turn, result in fast particle transport away from Earth due to enhanced outward radial diffusion driven by drift resonance of particles with ultralow frequency (ULF) waves. The loss into the interplanetary medium should be carefully separated from the precipitation into the atmosphere if any loss mechanism in the belts is investigated.

In this study, we present additional evidence for the ultrarelativistic electron loss in the heart of the outer belt that can be driven by resonant interaction with EMIC waves and complement the signatures of EMIC wave effects found by Usanova et al. (2014) for the time period from 9 October to 29 November 2012. We analyze phase space densities (PSD), derived from Van Allen Probe observations as a function of time and three adiabatic invariants  $\mu$ ,  $K$ , and  $L^*$  (Roederer, 2012; Schulz & Lanzerotti, 1974), to distinguish reversible (adiabatic) changes resulting from slow expansion and compression of the magnetic field from nonreversible (nonadiabatic) changes leading to particle loss or acceleration. The mechanisms causing nonreversible changes can be further differentiated if the PSD is considered as a function of  $L^*$  and time for constant  $\mu$  and  $K$  (Green & Kivelson, 2004; Reeves et al., 2013; Selesnick & Blake, 2000). In particular, local minimums in relativistic and ultrarelativistic electron PSD profiles are a distinctive feature of fast local loss, which can be driven by EMIC waves (Shprits et al., 2017). To provide observational evidence that EMIC waves scatter ultrarelativistic electrons into the Earth's atmosphere, we analyze such local minimums, narrowing of pitch angle distributions and the occurrence of EMIC waves on the ground.

## 2. Van Allen Probe Flux Measurements From 9 October to 29 November 2012

The interval of enhanced ultrarelativistic electron flux in the outer belt between 9 October and 29 November 2012 was teemed with multiple events of enhanced EMIC wave activity observed on the ground (Usanova et al., 2014) and gave us a good opportunity to gain insight into the processes driving the electron loss. Figure 1 illustrates the dynamics of differential ultrarelativistic electron fluxes from 9 October to 29 November. Note that the fluxes were averaged over the solid angle subtended by a whole sphere. Figure 1a shows the evolution of  $Dst$  and  $Kp$  indices, and Figures 1b–1d present the fluxes measured by the REPT instrument on board the Van Allen Probes. Figure 1e, adopted from Usanova et al. (2014), demonstrates the occurrence of EMIC waves on the ground measured by CARISMA stations at  $L = 4–4.5$ . The red circles in Figure 1e denote the days when EMIC waves were observed, while blue circles indicate the days when no EMIC wave activity was detected.



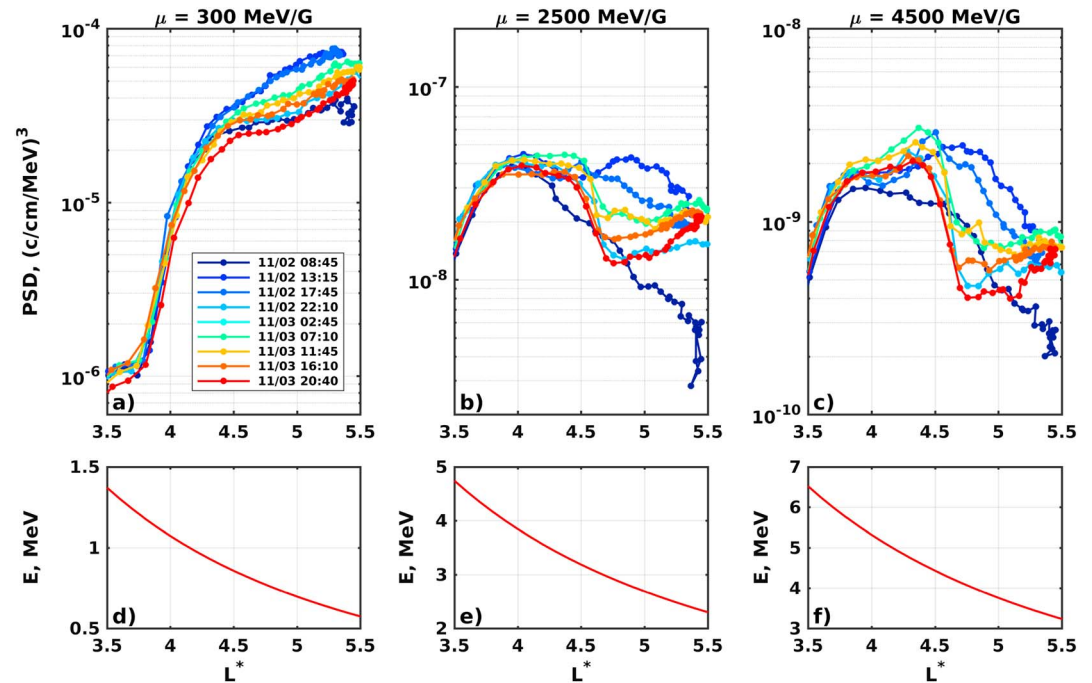
**Figure 1.** (a) *Kp* and *Dst* indices between 9 October and 29 November 2012, provided by OMNIWeb. (b–d) Differential electron fluxes measured by REPT on board Van Allen Probes in 2.6, 3.4, and 5.2 MeV energy channels as a function of time and *L\** computed using the TS07D magnetic field model (Tsyganenko & Sitnov, 2007). (e) The presence of EMIC waves measured by CARISMA stations at *L* = 4–4.5.

The prominent storm time flux dropouts are observed on 13 October, 1 November, and 13–14 November. The dropouts are accompanied by enhanced EMIC wave activity and minimums in the *Dst* index, which indicates that EMIC wave-induced scattering, adiabatic changes associated with the changes in the magnetic field configuration, magnetopause shadowing, and the loss into the interplanetary medium driven by outward radial diffusion can be responsible for such flux variations. Interestingly, EMIC waves are also observed not only during dropouts but also during the enhancement of the belt (e.g., 14–16 October) and intervals when the flux remains almost unchanged (e.g., 7 November 2012). It is therefore necessary to distinguish between competing loss and acceleration processes while performing analysis of EMIC wave effects on ultrarelativistic electrons in the belt.

### 3. Data and Methods

In order to extract signatures of EMIC wave-driven ultrarelativistic electron loss from in situ Van Allen Probe measurements, we applied the method proposed by Shprits et al. (2017). The method relies on the fact that interacting locally with electrons, EMIC waves produce pronounced local minimums in PSD profiles for constant first and second adiabatic invariants. Unless smoothed out by radial diffusion or local acceleration, the minimums can persist for hours and days, and they can be derived from the Van Allen Probe ECT suite observations.

We used 5 min averaged Magnetic Electron Ion Spectrometer (MagEIS) (Blake et al., 2013) and REPT data to calculate electron PSD as a function of adiabatic invariants. The first adiabatic invariant  $\mu$  has been calculated using local magnetic field observations measured by the Electric and Magnetic Field Instrument Suite and Integrated Science (Kletzing et al., 2013) on the satellites. To determine the invariants *K* and *L\** corresponding to the local pitch angles and position of satellites, we utilized the TS07D magnetic field model (Tsyganenko & Sitnov, 2007), which is implemented in the International Radiation Belt Environment Modeling (IRBEM) library (Boscher et al., 2012). We calculated PSD for  $K = 0.1G^{1/2}R_{Ei}$ , since this value of the second invariant roughly corresponds to 51°–55° pitch angles in the heart of the outer belt, and EMIC waves can efficiently resonate



**Figure 2.** (a–c) Evolution of PSD profiles derived from RBSP-A flux measurements as a function of  $L^*$  for  $K = 0.1G^{1/2} R_E$  and  $\mu = 300, 2,500$ , and  $4,500$  MeV/G on 2–3 November 2012 using the TS07D magnetic field model. The colors represent the end times of successive inbound and outbound satellite passes. (d–f) Energies corresponding to the chosen  $\mu$  and  $K$  calculated in the dipole field approximation.

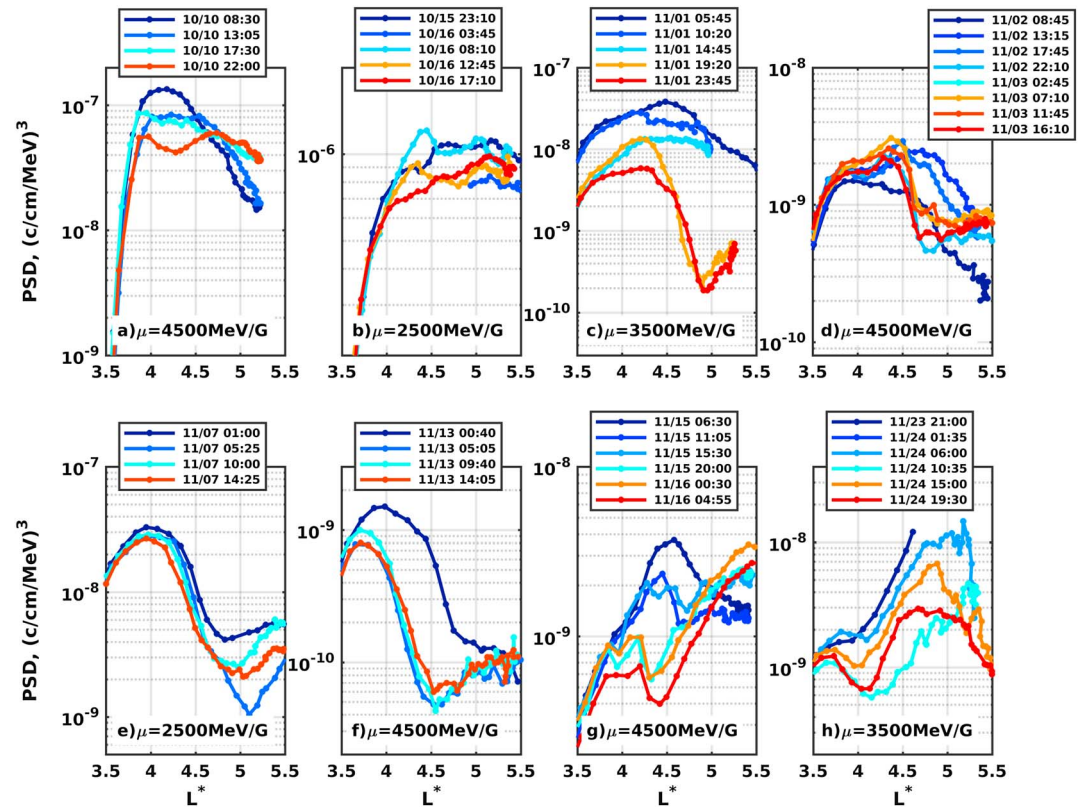
with the ultrarelativistic electrons of such pitch angles (Summers & Thorne, 2003). We used bilinear cubic interpolation to calculate PSD for particular values of  $\mu$  and  $K$  for constant  $L^*$ .

#### 4. Signatures of EMIC Wave-Driven Ultrarelativistic Electron Loss

Figure 2 illustrates the formation of local minimums on 2–3 November after the dropout of electron fluxes on 1 November (see Figures 1b–1d). Figures 2a–2c show the evolution of relativistic ( $\mu = 300$  MeV/G) and ultrarelativistic ( $\mu = 2,500$  and  $4,500$  MeV/G) PSD profiles. Figures 2d–2f depict energies corresponding to the chosen adiabatic invariants in the dipole field approximation. Predropout ultrarelativistic profiles, derived from RBSP-A measurements at around 08:45, 2 November, have wide pronounced peaks at  $L^* = 4$  and negative gradients above  $L^* \approx 4.3$  (see dark blue lines in Figures 2b and 2c). The next two passes of RBSP-A show an enhancement of both relativistic and ultrarelativistic electron PSD at higher L shells. The formation of local minimums can be noticed at the following satellite pass at  $L^* \approx 4.7$  around 22:10. The passes at  $\sim 07:10$  and  $\sim 11:45$  on 3 November show enhancement of ultrarelativistic electron PSD above  $L^* \approx 4.5$  without noticeable changes in the depth and width of the minimums. The subsequent passes demonstrate fast deepening of local minimums at  $L^* = 4.7$ , while profiles at  $L^* = 5.5$  do not show significant variations. Relativistic electron profiles, presented in Figure 2a, preserve monotonic behavior for the considered interval, showing only a slow gradual decrease.

To take into account uncertainties associated with calculations of  $K$  and  $L^*$  that globally depend on magnetic field configuration, we analyzed PSD profiles computed using T89 (Tsyganenko, 1989), T96 (Tsyganenko, 1995), T01S (Tsyganenko, 2002), and T04S (Tsyganenko & Sitnov, 2005) magnetic field models. Supporting information Figures S1–S4 correspond to the profiles derived on the basis of these models. The figures show that the local minimums in ultrarelativistic electron PSD profiles are observed independently of any known advanced magnetic field model, while relativistic electron profiles demonstrate a monotonic decrease with  $L^*$ .

The simultaneous formation of local minimums in ultrarelativistic electron PSD profiles at  $L^* = 4.7$ , the decrease between  $L^* = 4.5$  and  $5$ , and slight changes at  $L^* = 5.5$  imply a fast local loss process operating in a narrow region of L shells. The observed minimums could not be produced by magnetopause shadowing

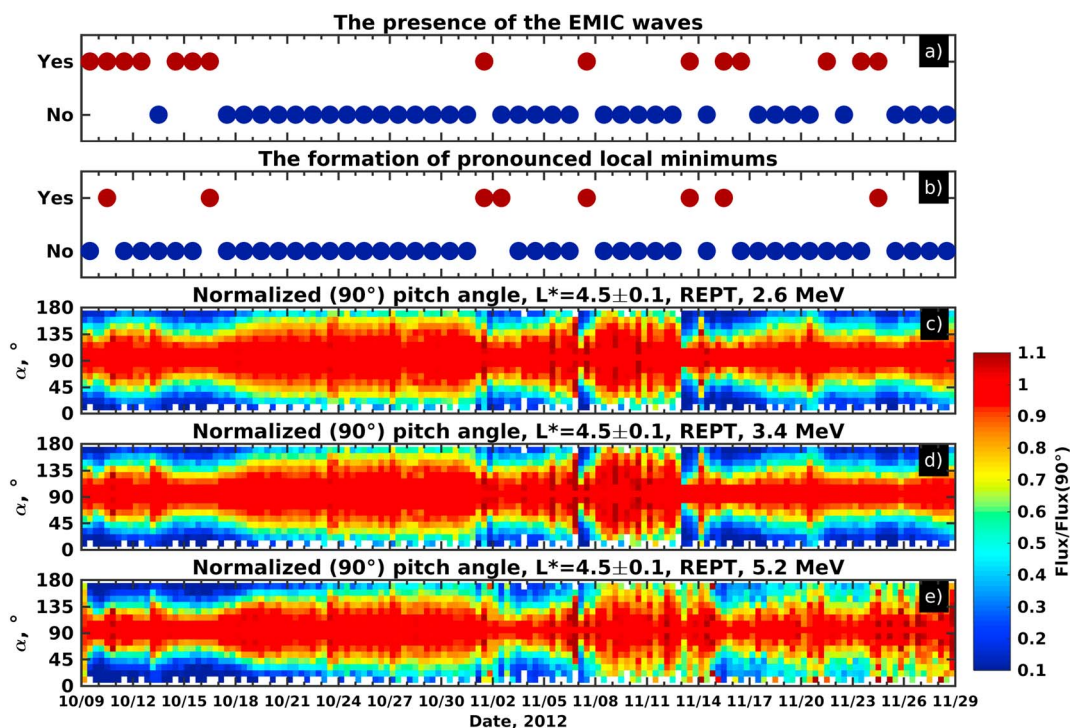


**Figure 3.** (a–h) Evolution of PSD profiles during periods when the most pronounced minimums were detected. The profiles are derived from RBSP-A flux measurements as a function of  $L^*$  for  $K = 0.1 \text{ G}^{1/2} R_E$  and constant  $\mu$  computed using the TS07D magnetic field model. Colors represent the end times of successive inbound and outbound satellite passes.

or by the local interaction with hiss or chorus waves, which is effective on the much longer time scales than the Van Allen Probe orbital period and characterized by a very weak dependence on radial distance (Orlova & Shprits, 2014; Orlova et al., 2014; Shprits et al., 2017). EMIC waves, resonating with ultrarelativistic electrons, can locally scatter electrons into the loss cone and form the minimums, as presented in Figures 2b and 2c. The absence of such minimums in the profiles of less energetic electrons (Figure 2a) is additional evidence of an EMIC wave-driven nature of loss, since less energetic electrons do not interact with EMIC waves according to the concept of minimum resonant energy (e.g., Summers & Thorne, 2003).

We analyzed PSD profiles for  $\mu = 2,500, 3,500,$  and  $4,500 \text{ MeV/G}$  from 9 October to 29 November 2016 and found eight events when noticeable local minimums were formed for at least one value of  $\mu$  in the chosen range. Figures 3a–3h show the profiles for the values of the first invariant  $\mu$  corresponding to the most pronounced minimums observed. The minimums differ in their position, depth, and width. The minimums are located between  $L^* = 4$  and  $5$ , and the width varies from  $\sim 0.3$  to  $\sim 1$  (in  $L$  shell units). The minimums appear simultaneously with the prominent dropouts (Figures 3c, 3f, and 3h) of PSD, or when no large variations of PSD detected (Figures 3a, 3b, 3d, 3e, and 3g). To diminish uncertainties associated with the magnetic field model, we also analyzed PSD profiles calculated using the T04S magnetic field model for identified events. Supporting information Figure S5 shows that the change in the magnetic field model mostly does not affect local minimums, though their shape can be slightly different.

Figure 4 summarizes the signatures of EMIC wave effects on ultrarelativistic electrons in the outer belt for the studied period. Figure 4a illustrates the occurrence of EMIC waves on the ground and is similar to Figure 1e. Figure 4b shows the presence of new pronounced local minimums in  $\mu = 2,500, 3,500,$  or  $4,500 \text{ MeV/G}$  profiles. Red markers (“Yes”) indicate days when the new local minimum was formed for at least one value in the chosen range, and blue markers (“No”) represent days without local minimums. Figures 4c–4e illustrate the pitch angle distributions of differential fluxes normalized by the equatorial flux at  $L^* = 4.5$  in the energy channels 2.6, 3.4, and 5.2 MeV.



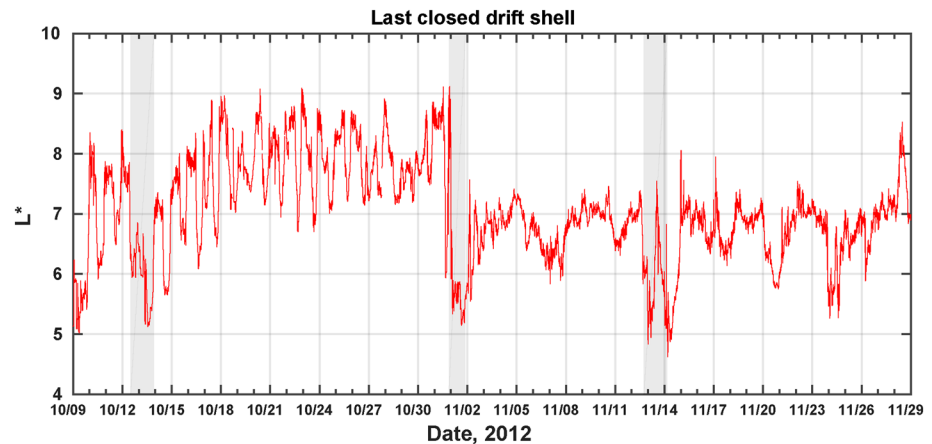
**Figure 4.** (a) The presence of EMIC waves measured by CARISMA stations between 9 October and 29 November 2012. (b) The formation of new local minima in PSD profiles for the same time period. (c–e) The pitch angle distributions of directional differential fluxes normalized by the equatorial flux as a function of time and pitch angles measured by Van Allen Probes in 2.6, 3.4, and 5.2 MeV energy channels at  $L^* = 4.5 \pm 0.1$ .

New local minima are formed during intervals of high EMIC wave activity, demonstrating strong correlation with the detection of EMIC waves on the ground (compare red dots in Figures 4a and 4b). The formation of local minima also coincides with the narrowing of pitch angle distributions, produced by the resonant interaction with EMIC waves (Usanova et al., 2014). The simultaneous formation of local minima, narrowing of the distributions, and EMIC wave detection on the ground is the observational evidence that EMIC waves not only change the shape of the normalized pitch angle distribution but may also result in particle loss.

## 5. Discussion

In this study, we analyzed 51 days of data measured by REPT on board the Van Allen Probes and found eight events of the formation of minima in ultrarelativistic electron PSD profiles. The events were accompanied by the narrowing of normalized pitch angle distributions and EMIC wave detection on the ground (see Figure 4). Such correlation shows from the statistical point of view that the new local minima have an EMIC wave-driven nature, as was found from the physical principles by Shprits et al. (2017). Complementing the findings of Usanova et al. (2014), which demonstrate that the narrowing of the pitch angle distributions coincides with EMIC wave measurements, the formation of new local minima may indicate EMIC wave-induced scattering of ultrarelativistic electrons.

The new local minima not only provide evidence of ultrarelativistic electron loss but can also help identify EMIC wave occurrence in space. Since the minima are the response of the electrons to the wave activity, they are the explicit manifestation of EMIC waves even if wave measurements are not available or spacecraft is not in the region of the waves. Therefore, the new minima can be used to support direct observations of the waves or to reveal EMIC waves that cannot be observed on the ground or in situ. Seven of eight local minima found between 9 October and 29 November 2012 serve as additional evidence of EMIC wave presence, while only one local minimum, which appeared on 2 November after the previous minimum on 1 November vanished, was not accompanied by EMIC waves and can be considered a potential indirect signature of the waves (see Figure 4).



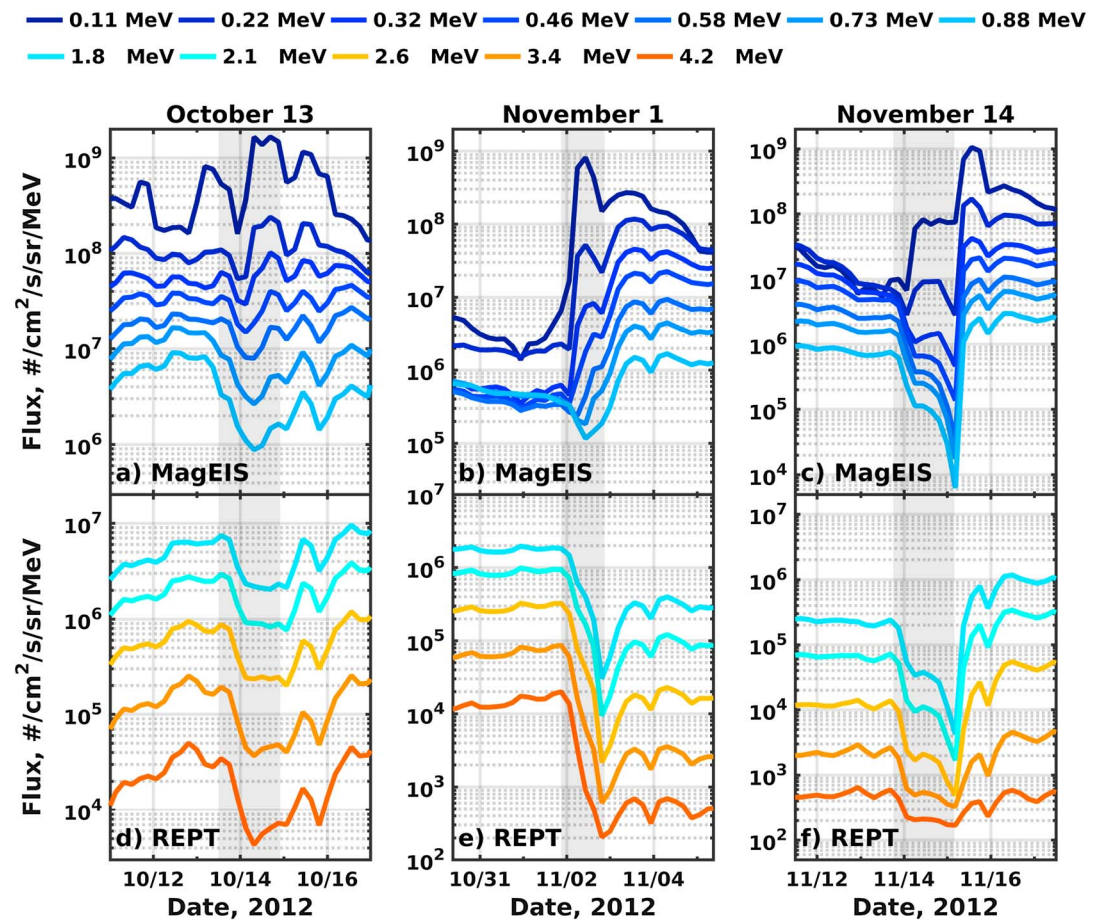
**Figure 5.** The last closed drift shell computed for equatorially mirroring particles using the TS07D magnetic field model. The gray shadowed regions correspond to significant storm time dropouts of ultrarelativistic electron fluxes at  $L^* = 5$ .

The duration of EMIC waves can be estimated from electron flux observations (converted to PSD) in case of deepening minimums as was discussed by Shprits et al. (2017). The deepening minimum in ultrarelativistic electron PSD profiles indicates that local loss induced by EMIC waves dominates over the acceleration process driven by radial diffusion that tends to smooth out the minimum. In this case, EMIC waves are present in the region of the deepening minimum at least as long as the deepening minimum is observed. However, if the minimum does not change its depth or the depth decreases, it is hard to distinguish EMIC wave effects on radiation belt electrons from the radial diffusion influence using only flux observations, and more detailed analysis involving EMIC wave parameters and radial diffusion rates is required. For this reason, we do not consider the question of EMIC wave duration in the current study and concentrate only on EMIC wave presence, which can be evident from the event of new local minimum formation.

The analysis of minimums can be complicated by different loss and acceleration processes that occur simultaneously with EMIC waves. For instance, Figure 3b shows the acceleration of particles after the minimum is formed at  $\sim 12:45$  on 16 October. In this case, the acceleration over the wide  $L^*$  range may accompany the loss of ultrarelativistic particles at  $L^* = \sim 4.6$ . Electron acceleration may render the minimums more shallow that is harder to detect visually or by using an automated algorithm.

For some events, EMIC waves were observed at times when local minimums are not seen in the data. That can be explained by the action of radial diffusion smoothing out the minimums, competition with local acceleration, or nonoptimal spectral properties of the waves for the efficient electron scattering. For instance, if the local minimum is formed at high  $L$  shells, it can be suddenly destroyed by magnetopause shadowing, as was detected on 12 November (see supporting information Figure S6). Acceleration due to inward radial diffusion or local interaction with nightside chorus waves can lead to the increase in PSD and smoothing out the minimums (see supporting information Figure S7 illustrating the disappearance of the minimums on 2 November). If competing processes significantly distort the minimums on time scales less than half of the Van Allen Probe period, the local minimums can be hardly observed, and therefore, their formation is detected more rarely than EMIC waves.

The local minimums and narrowing of pitch angle distributions are observed during the intervals of prominent storm time flux dropouts, indicating that EMIC waves can contribute to the noticeable flux decrease. It is necessary, however, to estimate the relative importance of EMIC wave-induced loss in the outer belt and the loss into the interplanetary medium driven by outward radial diffusion. Figure 5 shows the last closed drift shell calculated for equatorially mirroring electrons using the TS07D magnetic field model. To calculate the last closed drift shell, we find the farthest from the Earth magnetic field line that has only one local minimum at the noon magnetic longitude and calculate  $L^*$  corresponding to a given  $K$  value. This approach is implemented because of the limitations of IRBEM model computational framework to deal with double minimum field lines and provides an approximation of the last closed drift shell, which likely is at slightly higher  $L$  shells. A gray background color in Figure 5 marks the periods when the significant storm time dropouts of



**Figure 6.** Evolution of electron fluxes at  $L^* = 5 \pm 0.25$  measured by MagEIS and REPT instruments on board the RBSP-A during the most significant storm time dropouts that occurred on (a, d) 13 October, (b, e) 1 November, and (c, f) 13–14 November. The fluxes are averaged over the satellite orbital period. The gray shaded regions correspond to significant storm time dropouts of ultrarelativistic electron fluxes.

the ultrarelativistic electron fluxes were observed at  $L^* = 5$ . As can be clearly seen from the figure, all storm time events were accompanied by the sharp decrease of the last closed drift shell, which may indicate the enhanced electron loss into the interplanetary medium.

Figure 6 illustrates electron fluxes measured by the RBSP-A at  $L^* = 5$  and averaged over the satellite orbital period for the storms that occurred on 13 October, 1 November, and 13–14 November. The storms on 13 October and 13–14 November are characterized by a simultaneous dropout in electron fluxes from  $\sim 100$ – $200$  keV to several MeV energies (gray shaded regions in Figures 6a and 6d and Figures 6c and 6f), which can be indicative of the radiation belt variations associated with changes in magnetic field configuration. The event on 1 November, however, shows a significant decrease of multi-MeV electron fluxes and a simultaneous increase in  $\sim 100$ – $800$  keV fluxes (see Figures 6b and 6e). The flux enhancement can be driven by electron injections from the plasma sheet into the radiation belt region, concealing any possible effects of magnetopause shadowing on such electron population. Modeling studies (e.g., Drozdov et al., 2015; Kersten et al., 2014; Shprits et al., 2016) are required for further understanding of the balance between EMIC wave loss, loss into the interplanetary medium, and the electron source population injected at the nightside.

We emphasize the existence of essential coherence between EMIC wave observations and new local minimums in PSD profiles of ultrarelativistic electrons. Further studies are required for detailed understanding of correlations between the waves and the minimums, conditions under which the minimums are formed, and dependence of local minimum parameters on EMIC wave properties.



## Acknowledgments

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